temperature phase the static long-range ordered displacement wave (corresponding to the satellite reflections) takes the place of the thermal displacements in the high-temperature phase. The present results agree with the findings of inelastic neutronscattering studies (de Pater \& van Dijk, 1978; de Pater, Axe \& Currat, 1978). In these measurements a rather strong quasi-elastic diffuse peak was observed at high temperature ( 423 K ) which is centred in reciprocal space at a wave vector of $0.29 \mathbf{c}^{*}$. This diffuse scattering is due to short-range correlated displacement fluctuations and an energy analysis yielded a relaxation time for the fluctuations of $2 \times 10^{-11} \mathrm{~s}$, at 423 K .


Fig. 3. Temperature dependence of the isotropic thermal amplitude: $\left(\overline{u^{2}}\right)^{1 / 2} \equiv\left(B / 4 \pi^{2}\right)^{1 / 2}$, where $B$ is the overall Debye-Waller factor. The average splitting amplitude is also plotted, calculated from the data in Table 3.

In the present structure refinement these fluctuations in the normal high-temperature phase are seen as anisotropic smearing-out of the atoms.

In the low-temperature phase satellite reflections appear and the diffuse-scattering intensity decreases. The static displacements are seen as the splitting up of the atoms in the average structure. The splitting amplitude remains nearly constant in view of the decrease of the anisotropic thermal amplitude.

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# The Crystal Structure of $\mathbf{A g}_{2} \mathbf{I N O}_{3}$ 

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#### Abstract

$\mathrm{Ag}_{2} \mathrm{INO}_{3}$ is orthorhombic, space group $P 2_{1} 2_{1} 2_{1}$ with $a=7.4354$ (4), $b=7.7759$ (4), $c=17.1338$ (12) $\AA$, and $Z=8$. The structure was solved by direct methods from diffractometer data and was refined aniso-


tropically to an $R$ value of 0.029 . The I atoms are coordinated to six Ag atoms forming trigonal prisms which build up chains by sharing edges. The chains are linked together by either corners or edges. The $\mathrm{NO}_{3}$ groups are located in the cavities of this three-dimensional network.

## Introduction

This work is part of a systematic structural investigation of compounds of the general formula $\mathrm{Ag}_{n} X_{x} A_{y}$ ( $n>x$ ) with various $X, A, n$ and $x$. So far the detailed crystal structure is known for only three compounds meeting these stoichiometric requirements, viz $\mathrm{Ag}_{3} \mathrm{I}\left(\mathrm{NO}_{3}\right)_{2}$ (Birnstock \& Britton, 1970), $\mathrm{Ag}_{2} \mathrm{BrNO}_{3}$ (Persson \& Holmberg, 1977), and $\mathrm{Ag}_{13} \mathrm{I}_{9}\left(\mathrm{WO}_{4}\right)_{2}$ (Chan \& Geller, 1977) which is an ionic $\mathrm{Ag}^{+}$conductor. Brief reports on $\mathrm{Ag}_{2} \mathrm{IF} . \mathrm{H}_{2} \mathrm{O}$ (Holmberg \& Persson, 1975), $\mathrm{Ag}_{2} \mathrm{INO}_{3}$ and $\mathrm{Ag}_{7} \mathrm{I}_{2} \mathrm{~F}_{5} .2 \mathrm{H}_{2} \mathrm{O}$ (Persson, 1977) have been given.

## Experimental

Crystals of $\mathrm{Ag}_{2} \mathrm{INO}_{3}$ were obtained by saturation of an $\sim 1.5 \mathrm{M}$ aqueous solution of $\mathrm{AgNO}_{3}$ with AgI at about 363 K . After filtration at this temperature colourless rods crystallized on cooling. The crystals appeared to be specimens of either $\mathrm{Ag}_{2} \mathrm{INO}_{3}$ or $\mathrm{Ag}_{3} \mathrm{I}\left(\mathrm{NO}_{3}\right)_{2}$.
$\mathrm{Ag}_{2} \mathrm{INO}_{3}$ is stable in dry air at room temperature. The surface of the crystals darkens after exposure to light for a few days. This darkening had no serious effect on the X-ray intensities. No chemical analyses could be carried out since the crystals obtained were a mixture of the two compounds $\mathrm{Ag}_{2} \mathrm{INO}_{3}$ and $\mathrm{Ag}_{3} \mathrm{I}\left(\mathrm{NO}_{3}\right)_{2}$, and the total quantity of crystals was small. For the same reason the density was not determined experimentally.

Table 1 gives details of the crystal data, the collection of intensities, and the refinement. The method employed in data collection has been described (Elding, 1976). Weissenberg photographs revealed Laue class mmm and the systematic absences $h 00: h=2 n+1$, $0 k 0: k=2 n+1$ and $00 l: l=2 n+1$. A single-crystal diffractometer (CAD-4) was used for data collection. The cell dimensions were improved by least-squares refinement of 47 reflexions (Danielsson, Grenthe \& Oskarsson, 1976). The wavelength was $0.70930 \AA$.

The intensities of three standard reflexions $(13 \overline{5}, 248$, and $2,5, \overline{11}$ ), checked after every 100 measurements, showed a small decrease and the intensities were scaled with the expression $Y=1.0-\left(0.56 \times 10^{-3}\right) X$, where $X$ is the total time in hours that the crystal had been exposed to X-radiation. The values of $I$ and $\sigma_{c}(I)$, where $\sigma_{c}(I)$ is the standard deviation based on counting statistics, were corrected for Lorentz, polarization and absorption effects. The expression $p=\left(\cos ^{2} 2 \theta_{M}+\right.$ $\left.\cos ^{2} 2 \theta\right) /\left(1+\cos ^{2} 2 \theta_{M}\right)$ with $\theta_{M}=6.08^{\circ}$ was used in the correction for polarization. The crystal shape was described by ten planes.

## Structure determination and refinement

The heavy-atom structure was solved using the MULTAN system of computer programs (Germain, Main \& Woolfson, 1971). The O and N atoms were revealed in a subsequent difference synthesis. Fullmatrix least-squares refinement, minimizing $\sum_{w} w\left(F_{o} \mid-\right.$ $\left.\left|F_{c}\right|\right)^{2}$, was performed with weights $w=1 /\left[\sigma_{c}^{2} /\left(4 F_{o}^{2}\right)+\right.$ $\left.\left(a F_{o}\right)^{2}\right]$. The value of $a$ was chosen to give constant values of $\left\langle w\left(\left|F_{o}\right|-\left|F_{c}\right|\right)^{2}\right\rangle$ in different $\left|F_{o}\right|$ and $\sin \theta$ intervals. One scale factor, and positional and anisotropic thermal parameters were refined (Table 1).

Scattering factors were taken from International Tables for X-ray Crystallography (1974). The final refinement also included correction for extinction (Zachariasen, 1967) and anomalous dispersion by Ag and I. Only nine reflexions had extinction corrections $>10 \%$ in $\left|F_{o}\right|$.

In the last cycle the shifts in the parameters were less than $0.5 \%$ of the estimated standard deviations and the refinement was considered complete. A final difference synthesis showed peaks of height $1.35 \mathrm{e}_{\AA^{-3}}$ or less around the heavy atoms but it was otherwise featureless.

Fig. 1 shows a normal probability plot of $\delta R_{(i)}=$ $\left\lfloor\left|F_{o}(i)\right|-\left|F_{c}(i)\right|\right] / \sigma\left|F_{o}(i)\right|$ versus the values expected for a normal distribution (Abrahams \& Keve, 1971).

Table 1. Crystal data, the collection and reduction of intensities and the least-squares refinement

| $\mathrm{Ag}_{2} \mathrm{INO}_{3}: \mathrm{FW}$ | 404.65 |
| :---: | :---: |
| Crystal system | Orthorhombic |
| Space group | $P 2,2,2$, |
| $a(\AA)$ | 7.4354 (4) |
| $b$ ( $\AA$ ) | 7.7759 (4) |
| $c(\AA)$ | 17.1338(12) |
| $Z$ | 8 |
| $D_{x}\left(\mathrm{Mg} \mathrm{m}^{-3}\right)$ | 5.43 |
| Crystal size (mm) | $0.181 \times 0.056 \times 0.097$ |
| Radiation (graphite-monocromated) | Mo $K_{\text {a }}, \lambda=0.71073 \AA$ |
| Take-off angle ( ${ }^{\circ}$ ) | 3 |
| $\Delta \omega\left({ }^{\circ}\right)(\omega-2 \theta$ scan $)$ | $0.50+1.30 \tan \theta$ |
| $\Delta \theta\left({ }^{\circ}\right)$ | 3-30 |
| Minimum number of counts in a scan | 3000 |
| Maximum recording time (s) | 180 |


| $\mu\left(\right.$ Mo $\left.K_{1}\right)\left(\mathrm{mm}^{-1}\right)$ | 13.91 |
| :---: | :---: |
| Range of transmission factor | 0.364-0.565 |
| Number of measured reflexions | 1673 |
| Number of reflexions given zero weight ( $I<0$ ) | 193 |
| Number of independent reflexions used in the final refinement, $m$ | 1480 |
| Number of parameters refined, $n$ | 128 |
| $R=\Sigma\| \| F_{o^{\prime}}-{ }^{\prime} F_{c}\| \| / \Sigma\left\|F_{o}\right\|$ | 0.029 |
|  | 0.030 |
| $S=\left\|\Sigma w\left(\left\|F_{0}:-F_{c}\right\|\right)^{2} /(m-n)\right\|^{1 / 2}$ | 1.20 |
| $a$ (weighting function) | 0.015 |
| $g\left(\times 10^{-4}\right)$ (extinction) | 0.38 (2) |
| Mosaic spread (" of arc) | 15.4 |
| Domain size (mm) | $0.27 \times 10^{-3}$ |



Fig. 1. Normal probability plot of $1474 \delta R_{(i)}$ based on the structure factors.

Table 2. Positional parameters with estimated standard deviations

|  | $x$ | $y$ | $z$ |
| :--- | :--- | :--- | :--- |
| $\mathrm{Ag}(1)$ | $0.67435(11)$ | $0.03154(11)$ | $0.44214(4)$ |
| $\mathrm{Ag}(2)$ | $0.24780(11)$ | $0.02963(11)$ | $0.44081(4)$ |
| $\mathrm{Ag}(3)$ | $0.71726(11)$ | $0.53074(12)$ | $0.31602(5)$ |
| $\mathrm{Ag}(4)$ | $0.32554(11)$ | $0.48515(12)$ | $0.30493(4)$ |
| $\mathrm{I}(1)$ | $0.47074(8)$ | $0.84270(8)$ | $0.32144(3)$ |
| $\mathrm{I}(2)$ | $0.45908(8)$ | $0.35323(7)$ | $0.44995(3)$ |
| $\mathrm{O}(11)$ | $0.46911(96)$ | $0.44874(83)$ | $0.76430(31)$ |
| $\mathrm{O}(12)$ | $0.62920(94)$ | $0.30936(88)$ | $0.68016(42)$ |
| $\mathrm{O}(13)$ | $0.33518(98)$ | $0.30268(85)$ | $0.67359(36)$ |
| $\mathrm{O}(21)$ | $0.45349(99)$ | $0.67090(81)$ | $0.61412(34)$ |
| $\mathrm{O}(22)$ | $0.63834(100)$ | $0.76749(96)$ | $0.52591(38)$ |
| $\mathrm{O}(23)$ | $0.41976(103)$ | $0.93288(82)$ | $0.56900(34)$ |
| $\mathrm{N}(1)$ | $0.47594(94)$ | $0.35369(95)$ | $0.70566(36)$ |
| $\mathrm{N}(2)$ | $0.50243(96)$ | $0.79035(93)$ | $0.57119(37)$ |

The plot is close to linear with a slope of $1 \cdot 16$ and an intercept of 0.08 . The positive intercept indicates a systematic overestimation of $\left|F_{o}\right|-\left|F_{c}\right|$. This is because the method used in the measurements gives too small values for the background. The slope indicates that the values of $\sigma\left(F_{o}\right)$ are slightly underestimated. All computations were made on the Univac 1108 computer in Lund. Final positional parameters are given in Table 2.*

## Description of the structure

Selected interatomic distances and angles are given in Table 3. Fig. 2 gives a stereoscopic view of the contents of the unit cell. The structure consists of linked double

[^0]Table 3. Selected interatomic distances ( $\AA$ ) and angles $\left({ }^{\circ}\right)$ with estimated standard deviations

Symmetry code

| (i) $x$, - | -1+y, z | (ii) $\frac{1}{2}+x, \frac{1}{2}$ | $\frac{1}{2}-y, \quad 1-z$ |
| :---: | :---: | :---: | :---: |
| (iii) $-\frac{1}{2}+x$, | 位-y, 1-z | (iv) $1-x,-\frac{1}{2}$ | $\frac{1}{2}+y, \quad \frac{1}{2}-z$ |
| (v) $\frac{1}{2}+x$, | -1+y, $1-z$ | (vi) $\frac{3}{2}-x, 1$ | $1-y,-\frac{1}{2}+z$ |
| (vii) $\frac{1}{2}-x$, | $1-y,-\frac{1}{2}+z$ | (viii) $1-x, \frac{1}{2}$ | $\frac{1}{2}+y, \quad \frac{1}{2}-z$ |
| (ix) $\frac{1}{2}+x$, | 3 $\frac{3}{2}$ y, $1-z$ |  |  |
| $\mathrm{Ag}(1)-\mathrm{l}\left(1^{\prime}\right)$ | 2.954 (1) | $\mathrm{Ag}(3)-\mathrm{O}\left(11^{\text {vi }}\right.$ ) | 2.500 (7) |
| $\mathrm{Ag}(1)-\mathrm{I}\left(2^{\text {li }}\right.$ ) | 2.950 (1) | $\mathrm{Ag}(3)-\mathrm{O}\left(12^{\text {vi }}\right.$ ) | 2.875 (7) |
| $\mathrm{Ag}(1)-\mathrm{l}(2)$ | 2.973 (1) | $\mathrm{Ag}(3)-\mathrm{O}\left(13^{\text {ii) }}\right.$ ) | 2.743 (7) |
| $\mathrm{Ag}(2)-\mathrm{I}\left(1^{\prime}\right)$ | 3.007 (1) | $\mathrm{Ag}(3)-\mathrm{O}\left(21^{\text {lx }}\right.$ ) | $3 \cdot 147$ (7) |
| $\mathrm{Ag}(2)-\mathrm{I}(2)$ | 2.971 (1) | $\mathrm{Ag}(3)-\mathrm{O}\left(23^{\text {ix }}\right.$ ) | 2.496 (7) |
| $\mathrm{Ag}(2)-\mathrm{I}\left(2^{\text {iii) }}\right.$ ) | 2.990 (1) | $\mathrm{Ag}(4)-\mathrm{O}\left(11^{\text {vii) }}\right.$ ) | 2.356 (7) |
| $\mathrm{Ag}(3)-1(1)$ | 3.042 (1) | $\mathrm{Ag}(4)-\mathrm{O}\left(12^{\text {iiI }}\right.$ ) | 2.728 (7) |
| $\mathrm{Ag}(3)-\mathrm{l}\left(1^{\text {iv }}\right.$ ) | $3 \cdot 105$ (1) | $\mathrm{Ag}(4)-\mathrm{O}\left(13^{\text {vii) }}\right.$ ) | 3.036 (7) |
| $\mathrm{Ag}(3)-\mathrm{I}(2)$ | 3.295 (1) | $\mathrm{Ag}(1)-\mathrm{Ag}(2)$ | $3 \cdot 172$ (1) |
| $\mathrm{Ag}(4)-\mathrm{l}\left(1^{\text {iv }}\right.$ ) | 2.865 (1) | $\mathrm{Ag}(3)-\mathrm{Ag}(4)$ | 2.940 (1) |
| $\mathrm{Ag}(4)-\mathrm{I}(1)$ | 2.996 (1) | I(1)-I( $1^{\text {iv }}$ ) | 4.615 (1) |
| $\mathrm{Ag}(4)-\mathrm{l}(2)$ | 2.866 (1) | I(1)-l(1 $1^{\text {viil }}$ ) | 4.615 (1) |
| $\mathrm{Ag}(1)-\mathrm{O}\left(13^{1 \mathrm{i}}\right)$ | 2.650 (7) | I(2)-I(1) | 4.398 (1) |
| $\mathrm{Ag}(1)-\mathrm{O}\left(21^{1 \times}\right)$ | 2.778 (7) | I(2)-I(1) | 4.540 (1) |
| $\mathrm{Ag}(1)-\mathrm{O}\left(22^{\prime}\right)$ | 2.519 (7) | 1(2)-I( $1^{\text {iv }}$ ) | $4 \cdot 680$ (1) |
| $\mathrm{Ag}(1)-\mathrm{O}\left(23^{1}\right)$ | 2.983 (7) | I(2)-I(2ii) | 4.398 (1) |
| $\mathrm{Ag}(2)-\mathrm{O}\left(12^{\text {iii) }}\right.$ ) | 2.577 (7) | 1(2)-I( $2^{\text {ii }}$ ) | 4.398 (1) |
| $\mathrm{Ag}(2)-\mathrm{O}\left(22^{\text {iii) }}\right.$ ) | 2.847 (7) |  |  |
| $\mathrm{Ag}(2)-\mathrm{O}\left(22^{\text {iiii }}\right.$ ) | 2.515 (7) |  |  |
| $\mathrm{Ag}(2)-\mathrm{O}\left(23^{\prime}\right)$ | 2.651 (6) |  |  |
| $\mathrm{N}(1)-\mathrm{O}(11)$ | 1.248 (9) | $\mathrm{O}(11)-\mathrm{O}(12)$ | $2 \cdot 161$ (9) |
| $\mathrm{N}(1)-\mathrm{O}(12)$ | 1.268 (10) | $\mathrm{O}(11)-\mathrm{O}(13)$ | $2 \cdot 167$ (9) |
| $\mathrm{N}(1)-\mathrm{O}(13)$ | 1.247 (10) | $\mathrm{O}(12)-\mathrm{O}(13)$ | 2.190 (10) |
| $\mathrm{N}(2)-\mathrm{O}(21)$ | 1.239 (9) | $\mathrm{O}(21)-\mathrm{O}(22)$ | $2 \cdot 177$ (10) |
| $\mathrm{N}(2)-\mathrm{O}(22)$ | 1.286 (10) | $\mathrm{O}(21)-\mathrm{O}(23)$ | 2.193 (9) |
| $\mathrm{N}(2)-\mathrm{O}(23)$ | 1.268 (10) | $\mathrm{O}(22)-\mathrm{O}(23)$ | $2 \cdot 200$ (10) |
|  | $\mathrm{O}(11)-\mathrm{N}(1)-\mathrm{O}(12)$ | 118.35 (71) |  |
|  | $\mathrm{O}(11)-\mathrm{N}(1)-\mathrm{O}(13)$ | $120 \cdot 59$ (71) |  |
|  | $\mathrm{O}(12)-\mathrm{N}(1)-\mathrm{O}(13)$ | 121.06 (69) |  |
|  | $\mathrm{O}(21)-\mathrm{N}(2)-\mathrm{O}(22)$ | 119.01 (73) |  |
|  | $\mathrm{O}(21)-\mathrm{N}(2)-\mathrm{O}(23)$ | 122.03 (71) |  |
|  | $\mathrm{O}(22)-\mathrm{N}(2)-\mathrm{O}(23)$ | 118.93 (70) |  |
|  |  | 163.25 (4) |  |
|  | $\mathrm{I}\left(1^{1}\right)-\operatorname{Ag}(2)-\mathrm{I}\left(2^{\text {iii) }}\right.$ ) | 165.62 (4) |  |



Fig. 2. A stereoscopic pair of drawings showing the contents of the unit cell. Large circles are I and small Ag. Figs. 2, 3, and 5-8 have been drawn by the program ORTEP II (Johnson, 1971).
chains, along $\mathbf{b}$, containing trigonal-prismatic $\mathrm{Ag}_{6} \mathrm{I}$ units. The prisms share $\mathrm{Ag}-\mathrm{Ag}$ edges to form the double chains and these chains, two per unit cell, are linked by alternately shared $\mathrm{Ag}-\mathrm{Ag}$ edges and Ag corners. The chains are of a zigzag type. The $\mathrm{NO}_{3}$ groups reside in the cavities, formed by six Ag atoms, between the chains.
$\mathrm{AgINO}_{3}$ is a new structure type but can be related to the structure of tungsten carbide, WC, through a twinning operation and a translation of $\frac{1}{2} a$. In the WC structure type, every second $W$ trigonal prism is filled with C atoms. In the $\mathrm{Ag}_{2} \mathrm{INO}_{3}$ structure only every fourth of the Ag trigonal prisms is occupied. The twin plane is parallel with the $a b$ plane through the trigonal prisms which lie on different levels and share edges (Fig. 3). The twin blocks are two filled trigonal prisms wide and the structure can be written symbolically as 4 , 4, 4, 4 (Andersson \& Hyde, 1974).

Fig. 4(a) and (b) shows the $\mathrm{Ag}_{6} \mathrm{I}$ units of $\mathrm{I}(1)$ and $\mathrm{I}(2)$. The $\mathrm{Ag}-\mathrm{I}(1)$ and $\mathrm{Ag}-\mathrm{I}(2)$ distances range from 2.95-3.10 $\AA$ and $2 \cdot 86-2.99 \AA$ respectively, except for one, $\mathrm{Ag}(3)-\mathrm{I}(2)$, which is $3.30 \AA$. The distances are longer than those found in $\beta$-AgI (2.814 $\AA$; Burley, 1963) but of the same magnitude as those observed in $\mathrm{Ag}_{3} \mathrm{I}\left(\mathrm{NO}_{3}\right)_{2}(2 \cdot 84-3.08 \AA ;$ Birnstock \& Britton, 1970). In compounds $M_{y} \mathrm{Ag}_{n} \mathrm{I}_{x}$ with $x>n$ (Holmberg, 1976, references $26-39$ ) the $\mathrm{Ag}-\mathrm{I}$ distances fall in the approximate range $2 \cdot 80-3 \cdot 10 \AA$ irrespective of the nature of the counterion $M$; therefore no significant difference in the $\mathrm{Ag}-\mathrm{I}$ bonding strength can be inferred from such a comparison between bond lengths in the two general types of compounds.

The Ag surroundings are not easily described with a common polyhedron. Figs. 5-8 show the surroundings of the four different Ag atoms. $\mathrm{Ag}(1)$ has eight nearest neighbours: four O atoms, three I atoms, and one Ag atom. The O atoms and one of the I atoms make an irregular pentagon which approximately defines a plane. The five atomic positions deviate from the bestfitting least-squares plane by at most $0.07 \AA$ and $\mathrm{Ag}(1)$ is located only $0.002 \AA$ outside the plane. The other


Fig. 3. A projection of the structure along a. Large circles are I and small Ag. The marked coordinates are idealized (see Table 2).


Fig. 4. Stereoscopic drawings showing the trigonal prisms around (a) $\mathrm{I}(1)$ and (b) $\mathrm{I}(2)$.


Fig. 5. The environment of $\mathrm{Ag}(1)$.


Fig. 6. The environment of $\mathrm{Ag}(2)$.
two I atoms are situated at each side of this pentagon; the plane defined by $\mathrm{I}\left(1^{\mathrm{i}}\right), \mathrm{I}\left(2^{\mathrm{ii}}\right)$, and $\mathrm{Ag}(1)$ is close to perpendicular $\left(89.7^{\circ}\right)$ to the pentagon plane ( $c f$. the Ag environment in $\mathrm{Ag}_{2} \mathrm{BrNO}_{3}$ ). The distances are given in the figures. All $\mathrm{Ag}-\mathrm{O}$ distances in $\mathrm{Ag}_{2} \mathrm{INO}_{3}$ range from 2.36 to $3.04 \AA$ (Table 3), and they are in good agreement with $\mathrm{Ag}-\mathrm{O}$ distances found in $\mathrm{Ag}_{3} \mathrm{I}\left(\mathrm{NO}_{3}\right)_{2}$.
$\mathrm{Ag}(2)$ also has eight neighbouring atoms approximately in the same configuration as for $\mathrm{Ag}(1)$. The pentagon is somewhat less ideal than in the $\mathrm{Ag}(1)$ case. The deviations from the least-squares plane are at most $0.21 \AA, \mathrm{Ag}(2)$ being $0.056 \AA$ outside the plane. The $\mathrm{I}\left(1^{\mathrm{i}}\right), \mathrm{I}\left(2^{\mathrm{iii}}\right)$, and $\mathrm{Ag}(2)$ plane makes an angle of $84.9^{\circ}$ to the idealized plane of the five-membered ring. The position of the eighth atom, $\mathrm{Ag}(1)$ [and $\mathrm{Ag}(2)$ in the $\mathrm{Ag}(1)$ environment], is clear from Figs. 5 and 6. The polyhedra of $\mathrm{Ag}(1)$ and $\mathrm{Ag}(2)$ share an $\mathrm{O}(23)-\mathrm{I}(2)$ edge and an $I(1)$ corner.

The geometrical figure around $\mathrm{Ag}(3)$ can be described as a distorted monocapped trigonal anti-


Fig. 7. The environment of $\mathrm{Ag}(3)$.


Fig. 8. The environment of $\mathrm{Ag}(4)$.
prism (Fig. 7). An $\mathrm{Ag}(4)$ atom outside the I triangle makes the antiprism monocapped.

The $\mathrm{Ag}(4)$ environment is given in Fig. 8, showing the seven nearest neighbours. The polyhedra of $\mathrm{Ag}(3)$ and $\mathrm{Ag}(4)$ share a face via the triangle of I atoms.

The $\mathrm{NO}_{3}$ groups are located in the cavities between the chains of $\mathrm{Ag}_{6} \mathrm{I}$ trigonal prisms. These cation boxes (cf. $\mathrm{Ag}_{2} \mathrm{BrNO}_{3}$ ) are formed by six Ag atoms and they also have the shape of a trigonal prism. The $\mathrm{NO}_{3}$ group is oriented in this box with the O atoms directed towards the $\mathrm{Ag}-\mathrm{Ag}$ edges. The $\mathrm{NO}_{3}$ ions have close to perfect $D_{3 h}$ symmetry (Table 3).

Finally it should be observed that the $\mathrm{Ag}-X$ coordination in $\mathrm{Ag}_{2} X \mathrm{NO}_{3}$ is quite different for $X=\mathrm{Br}$ and $X=\mathrm{I}$. In the analogous bromide compound, Br is surrounded by five Ag atoms in a trigonal bipyramid. However, the structural similarities between $\mathrm{Ag}_{2} \mathrm{NO}_{3}$ and $\mathrm{Ag}_{3} \mathrm{I}\left(\mathrm{NO}_{3}\right)_{2}$ are quite striking, both compounds being built up by chains of trigonal $\mathrm{Ag}_{6} \mathrm{I}$ prisms.

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[^0]:    * Lists of structure factors, anisotropic thermal parameters and root-mean-square components of thermal displacement along the ellipsoid axes have been deposited with the British Library Lending Division as Supplementary Publication No. SUP 34032 (8 pp.). Copies may be obtained through The Executive Secretary, International Union of Crystallography, 5 Abbey Square, Chester CH 1 2HU, England.

